FreeFem++, a user language to solve PDE.

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http://www.freefem.org/ff++

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History

- 1987 MacFem/PCFem les ancêtres (O. Pironneau en Pascal) payant.
- 1992 FreeFem réécriture de C++ (P1,P0 un maillage) O. Pironneau, D. Bernardi, F. Hecht, C. Prudhomme (adaptation Maillage, bamg).
- 1996 FreeFem+ réécriture de C++ (P1,P0 plusieurs maillages) O. Pironneau, D. Bernardi, F. Hecht (algèbre de fonction).
- 1998 FreeFem++ réécriture avec autre noyau élément fini, et un autre langage utilisateur; F. Hecht, O. Pironneau, K.Ohtsuka.
- 1999 FreeFem 3d (S. Del Pino) , Une première version de freefem en 3d avec des méthodes de domaine fictif.
- 2008 FreeFem++ v3 réécriture du noyau élément fini pour prendre en compte les cas multidimensionnels : 1d,2d,3d...

For who, for what!

For what

- 0 R&D
- Academic Research ,
- 3 Teaching of FEM, PDE, Weak form and variational form
- 4 Algorithmes prototyping
- O Numerical experimentation
- **o** Scientific computing and Parallel computing

For who : the researcher, engineer, professor, student...

The mailing list mailto:Freefemppljll.math.upmc.fr with 414 members with a flux of 5-20 messages per day.

More than 2000 true Users (more than 100 download / month)

Main characteristics I/II

- Wide range of finite elements : continuous P1,P2 elements, discontinuous P0, P1, RT0,RT1,BDM1, elements ,Edge element, vectorial element, mini-element, ...
- Automatic interpolation of data from a mesh to an other one (with matrix construction if need), so a finite element function is view as a function of (x, y, z) or as an array.
- LU, Cholesky, Crout, CG, GMRES, UMFPack, SuperLU, MUMPS, HIPS, SUPERLU_DIST, PASTIX. ... sparse linear solver; eigenvalue and eigenvector computation with ARPACK.

- Automatic mesh generator, based on the Delaunay-Voronoï algorithm. (2d,3d (tetgen))
- Mesh adaptation based on metric, possibly anisotropic (only in 2d), with optional automatic computation of the metric from the Hessian of a solution. (2d,3d).
- Dynamic linking to add plugin.
- Full MPI interface
- Nonlinear Optimisation tools : CG, Ipopt, NLOpt, stochastic
- Wide range of examples : Navier-Stokes 3d, elasticity 3d, fluid structure, eigenvalue problem, Schwarz' domain decomposition algorithm, residual error indicator ...

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The C++ kernel / Dehli, (1992) (Idea, I)

```
My early step in C++
typedef double R;
class Cvirt { public: virtual R operator()(R) const =0;};
 class Cfonc : public Cvirt { public:
                                                                  // a function C
   R (*f)(R);
   R operator()(R x) const { return (*f)(x); }
  Cfonc( R (*ff)(R)) : f(ff) {} };
 class Coper : public Cvirt { public:
   const Cvirt *q, *d;
                                                               // the 2 functions
   R (*op)(R,R);
                                                                   // l'opération
   R operator()(R x) const { return (* \circ p)((*q)(x), (*d)(x));}
  Coper(R (*opp)(R,R), const Cvirt *qq, const Cvirt *dd):op(opp),q(qq),d(dd){}
   ~Coper() {delete q, delete d; } };
 static R Add(R a, R b) {return a+b;}
static R Sub(R a, R b) {return a-b;}
static R Mul(R a, R b) {return a*b;}
static R Div(R a, R b) {return a/b;}
 static R Pow(R a, R b) {return pow(a, b);}
```

A differential expression on in a PDE problem is like

 $f * [u_i | \partial_x u_i | \partial_y u_i | \dots] * [v_j | \partial_x v_j | \partial_y v_i | \dots]$

where [f, |g, ...] mean f or g, or ..., and where the unknown part is $[u_i|\partial_x u_i|\partial_y u_i|...] \equiv [(0,i)|(1,i)|(2,i)|...]$ is a pair of $i' \times i$, if we do the same of the test part, the differential expression is a formally sum of :

 $\sum_{k} f_k \times (i'_k, i_k, j'_k, j_k)$

So we can easily code this syntax :

```
varf a(u,v) = int2d(Th) (Grad(u)'*Grad(v)) -int2d(Th) (f*v)
+on(1,u=0);
matrix A=a(Vh,Vh,solver=UMFPACK);
real[int] b=a(0,Vh);
u[]=A^-1 * b;
```

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Laplace equation, weak form

Let a domain Ω with a partition of $\partial \Omega$ in Γ_2, Γ_e . Find u a solution in such that :

$$-\Delta u = 1 \text{ in } \Omega, \quad u = 2 \text{ on } \Gamma_2, \quad \frac{\partial u}{\partial \vec{n}} = 0 \text{ on } \Gamma_e$$
 (1)

Denote $V_g = \{v \in H^1(\Omega)/v_{|\Gamma_2} = g\}$. The Basic variational formulation with is : find $u \in V_2(\Omega)$, such that

$$\int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} 1v + \int_{\Gamma} \frac{\partial u}{\partial n} v, \quad \forall v \in V_0(\Omega)$$
(2)

The finite element method is just : replace V_g with a finite element space, and the FreeFem++ code :

Laplace equation in FreeFem++

The finite element method is just : replace V_g with a finite element space, and the FreeFem++ code :

```
mesh3 Th("fish3d.msh");// read a mesh 3d
fespace Vh(Th,P1); // define the P1 EF space
```

```
Vh u,v;// set test and unknow FE function in Vh.
macro Grad(u) [dx(u),dy(u),dz(u)] //EOM Grad def
solve laplace(u,v,solver=CG) =
    int3d(Th)( Grad(u)'*Grad(v))
    - int3d(Th)(1*v)
    + on(2,u=2); // int on γ2
plot(u,fill=1,wait=1,value=0,wait=1);
```

Run:fish.edp Run:fish3d.edp

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A first way to break complexity

Idea :

$$a(u,v) = \int_{\Omega} \nabla u.\nabla v$$

take an equi-partition of Ω in Ω_i for i=0 to N_p-1 the number of processor. then

$$a(u,v) = \sum_{i=0}^{N_p-1} \int_{\Omega_i} \nabla u . \nabla v$$

A first way to break complexity

Build matrix in parallel by assembling par region remark with the change function you change the region numbering to build region.
real c = mpisize/real(Th.nt);

Th=change(Th,fregion= min(mpisize-1,int(nuTriangle*c)));

2 Assemble the full matrix

③ Solve the linear using a good parallel solver (MUMPS)
load "MUMPS_FreeFem"
uh[] = A^-1*b;
// resolution

Run:Heat3d.edp

Run:NSCaraCyl-100-mpi.edp

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To solve the following Poisson problem on domain Ω with boundary Γ in $L^2(\Omega)$:

 $-\Delta u = f$, in Ω , and u = g on Γ ,

where $f \in L^2(\Omega)$ and $g \in H^{\frac{1}{2}}(\Gamma)$ are two given functions. Let introduce $(\pi_i)_{i=1,..,N_p}$ a positive regular partition of the unity of Ω , q-e-d :

$$\pi_i \in \mathcal{C}^0(\Omega): \quad \pi_i \ge 0 \text{ and } \sum_{i=1}^{N_p} \pi_i = 1.$$

Denote Ω_i the sub domain which is the support of π_i function and also denote Γ_i the boundary of Ω_i .

The parallel Schwarz method is Let $\ell = 0$ the iterator and a initial guest u^0 respecting the boundary condition (i.e. $u^0_{|\Gamma} = g$).

$$\forall i = 1.., N_p: \quad -\Delta u_i^{\ell} = f, \text{ in } \Omega_i, \quad \text{ and } u_i^{\ell} = u^{\ell} \text{ on } \Gamma_i$$

$$u^{\ell+1} = \sum_{i=1}^{N_p} \pi_i u_i^{\ell}$$
(4)

Some Remark

We never use finite element space associated to the full domain Ω because it to expensive. So we use on each domain i we defined $J_i = \{j \in 1, ..., N_p \mid \Omega_i \cap \Omega_j \neq \emptyset\}$ and we have

$$(u^{\ell+1})_{|\Omega_i} = \sum_{j \in J_i} (\pi_j u_j^{\ell})_{|\Omega_i}$$
(5)

We denote $u_{h|i}^{\ell}$ the restriction of u_{h}^{ℓ} on V_{hi} , so the discrete problem on Ω_{i} of problem (3) is find $u_{hi}^{\ell} \in V_{hi}$ such that :

$$\forall v_{hi} \in V_{0i} : \int_{\Omega_i} \nabla v_{hi} \cdot \nabla u_{hi}^\ell = \int_{\Omega_i} f v_{hi},$$

$$\forall k \in \mathcal{N}_{hi}^{\Gamma_i} : \sigma_i^k(u_{hi}^\ell) = \sigma_i^k(u_{h|i}^\ell)$$

where $\mathcal{N}_{hi}^{\Gamma_i}$ is the set of the degree of freedom (Dof) on $\partial\Omega_i$ and σ_i^k the Dof of V_{hi} .

Transfer Part equation(5)

```
To compute v_i = (\pi_i u_i)_{|\Omega_i} + \sum_{j \in J_i} (\pi_j u_j)_{|\Omega_i} and can be write the freefem++ function
Update with asynchronous send/recv (Otherwise dead lock).
func bool Update(real[int] &ui, real[int] &vi)
{ int n= jpart.n;
  for(int j=0;j<njpart;++j) Usend[j][]=sMj[j]*ui;</pre>
  mpiRequest[int] rq(n*2);
  for (int j=0; j<n;++j)</pre>
           Irecv(processor(jpart[j],comm,rq[j ]), Ri[j][]);
  for (int j=0; j<n;++j)</pre>
           Isend(processor(jpart[j], comm, rq[j+n]), Si[j][]);
  for (int j=0; j<n*2;++j)</pre>
           int k= mpiWaitAny(rq);
                                                           // set to (\pi_i u_i)_{|\Omega_i|}
  vi = Pii*ui;
                                  // apply the unity local partition .
   for(int j=0; j<njpart;++j)</pre>
                                                              // add (\pi_i u_i)_{|\Omega_i|}
      vi += rMj[j]*Vrecv[j][];
 return true; }
```

Finally you can easily accelerate the fixe point algorithm by using a parallel GMRES algorithm after the introduction the following affine S_i operator sub domain Ω_i .

Where the parallel MPIGMRES or MPICG algorithm is to solve $A_i x_i = b_i, i = 1, ..., N_p$ by just changing the dot product by reduce the local dot product of all process with the following MPI code :

```
A simple coarse grid is we solve the problem on the coarse grid :
func bool CoarseSolve(real[int]& V, real[int]& U,
             mpiComm& comm)
    if (AC.n==0 && mpiRank(comm)==0)
                                               // first time build
      AC = vPbC(VhC,VhC,solver=sparsesolver);
   real[int] Uc(Rci.n), Bc(Uc.n);
   Uc= Rci*U:
                                                 // Fine to Coarse
   mpiReduce(Uc, Bc, processor(0, comm), mpiSUM);
   if (mpiRank(comm) == 0)
      Uc = AC^{-1}*Bc;
                                                // solve of proc 0
    broadcast (processor(0, comm), Uc);
   V = Pci * Uc;
                                                 11
                                                        Coarse to Fine
```

Limitation : if the initial problem, data have oscillation, you must use homogenization technic on coarse problem, or use the F. Nataf and co, preconditionner.

So we finally we get 4 algorithms

- **1** The basic schwarz algorithm $u^{\ell+1} = S(u^{\ell})$, where S is one iteration of schwarz process.
- 2 Use the GMRES to find u solution of the linear system Su u = 0.
- **3** Use the GMRES to solve parallel problem $A_i u_i = b_i$, $i = 1, ..., N_p$, with RAS precondicionneur
- **4** Use the method with two level precondicionneur RAS and Coarse.

On the SGI UV 100 of the lab :



A Parallel Numerical experiment on laptop

We consider first example in an academic situation to solve Poisson Problem on the cube $\Omega=]0,1[^3$

 $-\Delta u = 1, \text{ in } \Omega; \qquad u = 0, \text{ on } \partial \Omega.$ (6)

With a cartesian meshes \mathcal{T}_{hn} of Ω with $6n^3$ tetrahedron, the coarse mesh is \mathcal{T}_{hm}^* , and m is a divisor of n.

We do the validation of the algorithm on a Laptop Intel Core i7 with 4 core at 1.8 Ghz with 4Go of RAM DDR3 at 1067 Mhz,

Run:DDM-Schwarz-Lap-2dd.edp Run:DDM-Schwarz-Lame-3d.edp Run:DDM-Schwarz-Lame-2d.edp Run:DDM-Schwarz-Stokes-2d.edp

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To solve F(u) = 0 the Newton's algorithm is

• u^0 a initial guest

🙆 do

- find w^n solution of $DF(u^n)w^n = F(u^n)$
- $u^{n+1} = u^n w^n$
- if($||w^n|| < \varepsilon$) break;

For Navier Stokes problem the algorithm is : $\forall v, q$,

$$F(u,p) = \int_{\Omega} (u.\nabla)u.v + \nu\nabla u : \nabla v - q\nabla . u - p\nabla . v + BC$$

$$DF(u,p)(w,w_p) = \int_{\Omega} (w.\nabla)u.v + (u.\nabla)w.v + \int_{\Omega} \nu\nabla w : \nabla v - q\nabla .w - w_p\nabla .v + BC0$$

Run:cavityNewtow.edp Run:Hyper-Elasticity-2d.edp Run:NSNewtonCyl-100-mpi.edp

In Euclidean geometry the length $|\gamma|$ of a curve γ of \mathbb{R}^d parametrized by $\gamma(t)_{t=0..1}$ is

$$\gamma| = \int_0^1 \sqrt{<\gamma'(t),\gamma'(t)>} dt$$

We introduce the metric $\mathcal{M}(x)$ as a field of $d \times d$ symmetric positive definite matrices, and the length ℓ of Γ w.r.t \mathcal{M} is :

$$\ell = \int_0^1 \sqrt{<\gamma'(t), \mathcal{M}(\gamma(t))\gamma'(t)>} dt$$

The key-idea is to construct a mesh where the lengths of the edges are close to 1 accordingly to $\mathcal{M}.$

The mains IDEA for mesh generation

- The difficulty is to find a tradeoff between the error estimate and the mesh generation, because this two work are strongly different.
- $\bullet\,$ To do that, we propose way based on a metric ${\cal M}$ and unit mesh w.r.t ${\cal M}$
- The metric is a way to control the mesh size.
- remark : The class of the mesh which can be created by the metric, is very large.

Idea : Metrix intersection The unit ball $\mathcal{B}_{\mathcal{M}}$ in a metric \mathcal{M} plot the maximum mesh size on all the direction, is a ellipse. If you we have two unknowns u and v, we just compute the metric \mathcal{M}_u and \mathcal{M}_v , find a metric \mathcal{M}_{uv} call intersection with the biggest



ellipse such that : $\mathcal{B}_{\mathcal{M}_{uv}} \subset \mathcal{B}_{\mathcal{M}_u} \cap \mathcal{B}_{\mathcal{M}_v}$

Build of the metric form the solution u

Optimal metric norm for interpolation error (function adaptmesh in freefem++) for P_1 continuous Lagrange finite element

- $L^{\infty}: \mathcal{M} = \frac{1}{\varepsilon} |\nabla \nabla u| = \frac{1}{\varepsilon} |\mathcal{H}|$ where $\mathcal{H} = \nabla \nabla u$
- $L^p: \mathcal{M} = \frac{1}{\varepsilon} |det(\mathcal{H})|^{\frac{1}{2p+2}} |\mathcal{H}|$ (result of F. Alauzet, A. Dervieux)

In Norm $W^{1,p}$, the optimal metric \mathcal{M}_{ℓ} for the P_{ℓ} Lagrange finite element, Optimal is given by (with only acute triangle) (thank J-M. Mirebeau)

$$\mathcal{M}_{\ell,p} = \frac{1}{\varepsilon} (det \mathcal{M}_{\ell})^{\frac{1}{\ell_p+2}} \mathcal{M}_{\ell}$$

.

and (see MetricPk plugin and function)

• for $P_1 : \mathcal{M}_1 = \mathcal{H}^2$ (sub optimal with acute triangle take \mathcal{H})

• for
$$P_2: \mathcal{M}_2 = 3\sqrt{\begin{pmatrix} a & b \\ b & c \end{pmatrix}^2 + \begin{pmatrix} b & c \\ c & a \end{pmatrix}^2}$$
 with
 $D^{(3)}u(x,y) = (ax^3 + 3bx^2y + 3cxy^2 + dy^3)/3!$

Example of adaptation process

Find optimal mesh in norm L^∞ to represent :

$$u = (10 * x^{3} + y^{3}) + atan2(0.001, (sin(5 * y) - 2 * x))$$
$$v = (10 * y^{3} + x^{3}) + atan2(0.01, (sin(5 * x) - 2 * y)).$$



The IPOPT optimizer in a FreeFem++ script is done with the <code>IPOPT</code> function included in the <code>ff-Ipopt</code> dynamic library.

IPOPT is designed to solve constrained minimization problem in the form :

 $\begin{array}{ll} \text{find} & x_0 = \mathop{\mathrm{argmin}}_{x \in \mathbb{R}^n} f(x) \\ \text{s.t.} & \left\{ \begin{array}{ll} \forall i \leq n, \; x_i^{\text{lb}} \leq x_i \leq x_i^{\text{ub}} & (\text{simple bounds}) \\ \forall i \leq m, \; c_i^{\text{lb}} \leq c_i(x) \leq c_i^{\text{ub}} & (\text{constraints functions}) \end{array} \right. \end{array}$

Where ub and lb stand for "upper bound" and "lower bound". If for some $i, 1 \leq i \leq m$ we have $c_i^{\text{lb}} = c_i^{\text{ub}}$, it means that c_i is an equality constraint, and an inequality one if $c_i^{\text{lb}} < c_i^{\text{ub}}$.

Bose Einstein Condensate

Just a direct use of Ipopt interface (2 day of works) The problem is find a complex field u on domain \mathcal{D} such that :

$$u = \underset{||u||=1}{\operatorname{argmin}} \int_{\mathcal{D}} \frac{1}{2} |\nabla u|^2 + V_{trap} |u|^2 + \frac{g}{2} |u|^4 - \Omega i \overline{u} \left(\left(\frac{-y}{x} \right) \cdot \nabla \right) u$$

to code that in $\ensuremath{\mathsf{FreeFem}}\xspace+$

use

- lpopt interface (https://projects.coin-or.org/Ipopt)
- Adaptation de maillage

Run:BEC.edp

Solid to Liquid and Natural Convection

The starting point is almost the Orange Problem is describe in web page http://www.ljll.math.upmc.fr/~hecht/ftp/ff++days/2011/Orange-problem.pdf. The coupling of natural convection modeled by the Boussinesq approximation and liquid to solid phase change in $\Omega =]0, 1[^2$, No slip condition for the fluid are applied at the boundary and adiabatic condition on upper and lower boundary and given temperature θ_r (resp θ_l) at the right and left boundaries.

The model is : find the field : the velocity $\boldsymbol{u} = (u_1, u_2)$, the pressure p and temperature θ :

$$\begin{array}{cccc} \boldsymbol{u} & \text{given} & \text{in } \Omega_s \\ \partial_t \boldsymbol{u} + (\boldsymbol{u} \nabla) \boldsymbol{u} + \nabla . \mu \nabla \boldsymbol{u} + \nabla p &= -c_T \boldsymbol{e}_2 & \text{in } \Omega_f \\ \nabla . \boldsymbol{u} &= 0 & \text{in } \Omega_f \\ \partial_t \theta + (\boldsymbol{u} \nabla) \theta + \nabla . k_T \nabla \theta &= \partial_t S(T) & \text{in } \Omega \end{array}$$

Where Ω_f is the fluid domain and the solid domain is $\Omega_s = \Omega \setminus \Omega_f$.

(7)

Solid to Liquid and Natural Convection

The enthalpy of the change of phase is given by the function S; μ is the relative viscosity, k_T the thermal diffusivity.

In $\Omega_f = \{x \in \Omega; \theta > \theta_f\}$, with θ_m the melting temperature the solid has melt. We modeled, the solid phase as a fluid with huge viscosity, so :

$$\mu = \left\{ \begin{array}{ccc} \theta < \theta_f & \sim & 10^6 \\ \theta \ge \theta_m & \sim & \frac{1}{\mathrm{Re}} \end{array} \right. ,$$

The Stefan enthalpy S_c with defined by $S_c(\theta) = H(\theta)/S_{th}$ where S_{the} is the stefan number, and H is the Heaviside function with use the following smooth the enthalpy :

$$S(\theta) = \frac{\tanh(50(\theta - \theta_m)))}{2S_{te}}$$

The true device



the Algorithm

We apply a fixed point algorithm for the phase change part (the domain Ω_f is fixed at each iteration) and a full no-linear Euler implicit scheme with a fixed domain for the rest. We use a Newton method to solve the non-linearity.

- if we don't make mesh adaptation, the Newton method do not converge
- if we use explicit method diverge too,
- $\bullet\,$ if we implicit the dependance in Ω_s the method also diverge.

This is a really difficult problem.

the Algorithm, implementation

The finite element space to approximate $u1, u2, p, \theta$ is defined by

```
fespace Wh(Th,[P2,P2,P1,P1]);
```

We do mesh adaptation a each time step, with the following code :

This mean, we adapt with all variable plus the 2 melting phase a time n + 1 and n and we smooth the metric with a ratio of 1.2 to account for the movement of the melting front.

The Newton loop

the fixed point are implemented as follows

```
real err=1e100,errp ;
for(int kk=0;kk<2;++kk)//2 step of fixe point on Ω<sub>s</sub>
{ nu = nuT; // recompute the viscosity in Ω<sub>s</sub>,Ω<sub>f</sub>
  for(int niter=0;niter<20; ++niter)//Newton loop
  { BoussinesqNL;
    err = u1w[].linfty;
    cout << niter << "_err_NL_" << err <<endl;
    u1[] -= u1w[];
    if(err < tolNewton) break; }// convergence
}</pre>
```

The linearized problem

```
problem BoussinesqNL([u1w,u2w,pw,Tw],[v1,v2,q,TT])
= int2d(Th) ( [u1w, u2w, Tw]' * [v1, v2, TT] * cdt
     + UgradV(u1,u2,u1w,u2w,Tw)' * [v1,v2,TT]
     + UgradV(u1w,u2w,u1,u2,T)' * [v1,v2,TT]
     + ( Grad(u1w,u2w)'*Grad(v1,v2)) * nu
     + ( Grad(u1,u2)'*Grad(v1,v2)) * dnu* Tw
     + cmT*Tw*v2 + grad(Tw)'*grad(TT)*kT
     - \operatorname{div}(u1w, u2w) \star q - \operatorname{div}(v1, v2) \star pw - eps \star pw \star q
     + dS(T) *Tw*TT*cdt )
   - int2d(Th)(
      [u1,u2,T]'*[v1,v2,TT]*cdt
     + UgradV(u1,u2,u1,u2,T)' * [v1,v2,TT]
     + ( Grad(u1,u2)'*Grad(v1,v2)) * nu
     + cmT*T*v2 - eps*p*q + grad(T)'*grad(TT)*kT
     - div(u1,u2)*g -div(v1,v2)*p
     + S(T) *TT * cdt - [u1p, u2p, Tp]' * [v1, v2, TT] * cdt

    S(Tp)*cdt*TT)

 + on (1, 2, 3, 4, u1w=0, u2w=0) + on (2, Tw=0) + on (4, Tw=0);
```

The parameters of the computation

take case 2 from

Shimin Wang, Amir Faghri, and Theodore L. Bergman. A comprehensive numerical model for melting with natural convection. *International Journal of Heat and Mass Transfer*, January 2010.

 $\theta_m = 0$, Re = 1, $S_{te} = 0.045$, $P_r = 56.2$, $R_a = 3.27 \ 10^5$, $\theta_l = 1$, $\theta_r = -0.1$ so in this case cmT = $c_T = -R_a/P_r$, kT = $k_T = 1/P_r$, eps = 10^{-6} , time step $\delta t = 10^{-1}$, cdt = $1/\delta t$, at time t = 80 and we get a good agreement with the article.

Variables and a-dimension

The fusion temperature will be denoted by T_f . Using a lengthscale $L_{ref} = H$ and a liquid reference state $(\rho_{ref}, V_{ref}, T_{ref})$, we can define the following scaling for the space, velocity, temperature and time variables :

$$\vec{x} = \frac{\vec{X}}{L_{ref}}, \quad \vec{u} = \frac{\vec{U}}{V_{ref}}, \quad \theta = \frac{T - T_{ref}}{T_h - T_c}, \quad t = \frac{\tau}{t_{ref}}, \quad t_{ref} = L_{ref}/V_{ref}.$$

 $f_B(\theta) = \frac{\mathcal{R}a}{\mathcal{P}r \mathcal{R}e^2} \theta$, The buoyancy force (8)

where the Reynolds, Prandtl and Rayleigh numbers, are defined as :

$$\mathcal{R}e = \frac{\rho_{ref}V_{ref}L_{ref}}{\mu_l}, \quad \mathcal{P}r = \frac{\nu_l}{\alpha_l}, \quad \mathcal{R}a = \frac{g\beta L_{ref}^3(T_h - T_c)}{\nu_l\alpha_l}, \tag{9}$$

with μ denoting the viscosity, ν the kinematic viscosity, α the thermal diffusivity, β the thermal expansion coefficient and g the gravitational acceleration.

Phase change with Natural Convection





So now, a real problem, get the physical parameter of the real experiment. Run:Orange-Newton.edp

Density of water

Pure water exhibits a nonlinear density variation for $T < 10.2 \,^{\circ}\text{C}$ with a maximum at $T_m = 4.0293 \,^{\circ}\text{C}$. We use below the following density-temperature relationship :

$$\rho(T) = \rho_m \left(1 - w \left| T - T_m \right|^q \right), \tag{10}$$

with $\rho_m = 999.972 \text{ [kg/m}^3$], $w = 9.2793 \cdot 10^{-6} \text{ [(}^{\circ}C)^{-q}\text{]}$, and q = 1.894816. Choosing the fusion temperature $T_f = 0 \,^{\circ}\text{C}$ as reference, the bouyancy term $f_B = g(\rho_{ref} - \rho)/\rho_{ref}$ appearing becomes after scaling :

$$f_B(\theta) = \frac{\mathcal{R}a}{\mathcal{P}r\,\mathcal{R}e^2} \frac{1}{\beta(T_h - T_c)} \frac{\rho(\theta_f) - \rho(\theta)}{\rho(\theta_f)},\tag{11}$$

where $\beta = (1/\rho_m) (d\rho/dT)$ is the thermal expansion coefficient with the value $\beta = 6.91 \cdot 10^{-5}$ [(K)⁻¹]. Note that (11) compared to the classical linear form (8) non only introduces a new nonlinear term, but also the coefficient in front of this term is very large, since proportional to $\mathcal{R}a/(\beta(T_h - T_c))$.



Natural convection of water in a differentially heated cavity. Problem definition and streamlines of the steady flow.

Movie:mesh

The 9th Workshop on Petascale Computing, 13/06/2013.

Movietlow

Movie:flows Region

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Freezing of pure water



Freezing of pure water : Configuration at (physical time) $\tau = 2340[s]$: (a) experimental image; the thick red line represents the solid-liquid interface computed with the present method (b) Movie:flow Movie:mesh Movie:flows Region

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Freezing of pure water



Freezing of pure water. Computed configuration at (physical time) $\tau = 2340[s]$: (a) finite-element mesh refined along the solid-liquid interface ($T = 0 \,^{\circ}$ C) and also along the line of maximum water density ($T = 4 \,^{\circ}$ C) (b) temperature iso-lines.

Introduction

- 2 Academic Examples
- 3 A first way to break complexity
- 4 Schwarz method with overlap
- **5** Some tools and examples
- 6 Future/Conclusion

Freefem++ v3 is

- \bullet very good tool to solve non standard PDE in 2D/3D
- to try new domain decomposition domain algorithm

The the future we try to do :

- Build more graphic with VTK, paraview , ... (in progress)
- Add Finite volume facility for hyperbolic PDE (just begin C.F. FreeVol Projet)
- 3d anisotrope mesh adaptation
- automate the parallel tool

Thank for you attention.